

Energy Partitioning and Microstructural Observations Related to Perforation of Titanium and Steel Targets

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Abstract

This report presents an analysis of two target materials and the associated energetics related to the initial penetration into the target and perforation as the penetrator exits the target. Impact tests were conducted for tungsten alloy rods striking rolled homogeneous armor (RHA) and titanium alloy plates. Rod impact velocities were nominal 1,500 and 2,000 m/s. Target thicknesses were chosen so that the rods would overmatch the targets and lose some 200 m/s during penetration. The tests utilized flash x-rays to determine rod residual lengths and velocities, and target plug features, to include thicknesses and velocities. From these observables, experimental determination of the corresponding kinetic energies (KEs) and estimates for the fracture energies were obtained. Also, in each case, target material adjacent to penetration channel walls was examined by optical and electron microscopy and x-ray diffraction to gain further insight into deformation processes (cavity expansion) during penetration. The analytic penetration model gave results that were in good agreement with the experimental observables. In addition, it was observed that the RHA follows traditional plastic flow of cavity expansion while titanium alloy shows deformation features that deviate significantly. The report discusses possible causes for these differences.

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1. Introduction

Recent advances in theoretical penetration mechanics by Grace [1] allow the problem of kinetic energy (KE) rod impact with plates at impact velocities extending to hypervelocity to be separated into two phases. These phases include (1) an initial penetration into the target and (2) a following perforation of the target as the rod exits the target rear surface. The model used in this study provided a framework through which the overall energetics could be analyzed theoretically as demonstrated by Rupert and Grace [2]. Previous analytic approaches to the problem have included adaptations of the Tate analysis for semi-infinite targets by Zook [3], the semi-empirical formulas of Lambert [4], and applications of "hydrocodes" [1, 3, 5], for example. Collectively, these approaches have produced a large body of knowledge about penetration and perforation that allows engineering estimates of target efficiencies against long rods. However, certain aspects of the penetration process have not been adequately addressed. These aspects include the nature of target deformation in terms of changes in microstructure of the material, where and how the KE of the rod is dissipated in the target, and details of target material failure and fracture during perforation.

A recently developed model for penetration/perforation of single plates [1] is used to track details of the penetration events at impact velocities extending to the hydrodynamic regime, wherein the process is considered to take place in two separate phases. These phases include (1) an initial penetration into the plate target and (2) a following perforation of the target as the rod exits the target rear surface. Since the analysis provides mass-velocity relationships and estimates for the fracture energies, it is used here to analyze the energetics of penetration and perforation processes. The present work also includes experimental test results for tungsten alloy (WA) long rods that were fired at high velocity against two targets of practical interest (i.e., rolled homogeneous armor [RHA] and titanium-6% aluminum-4% vanadium alloy [Ti-6/4]). Optical and transmission electron microscopy examinations were conducted on the target materials after perforation. X-ray diffraction was conducted to determine any structural phase shifts in the Ti-6/4 material. This report includes results that demonstrate significant differences in the impact response of RHA and Ti-6/4 in terms of the

dissipation of energy in these targets and presents supporting calculations and experimental observations.

2. Experimental Approaches and Results

2.1 Rod-Impact Tests. Rod-impact tests were designed so that the penetration efficiencies of RHA and Ti-6/4 plate targets could be compared. The penetrators were fired from a laboratory gun consisting of a Bofors' 40-mm gun breech assembly with a custom-made 40-mm smoothbore barrel that was positioned approximately 3 m in front of the targets. High-speed (flash) radiography was used to record and measure projectile pitch, yaw, striking velocity, residual rod length, residual rod velocity, and plug thickness and velocity. Two pairs of orthogonal x-ray tubes were positioned in the vertical and horizonal planes along the shot line in front of the target (as illustrated in Figure 1). An additional pair of x-ray tubes was positioned in the horizonal plane along the shot line behind the target. Propellant weight was adjusted for desired nominal velocities (v_s) of 1,500 m/s and 2,000-m/s. Shots that had a total yaw in excess of 2° were considered "no tests," and those data were disregarded.

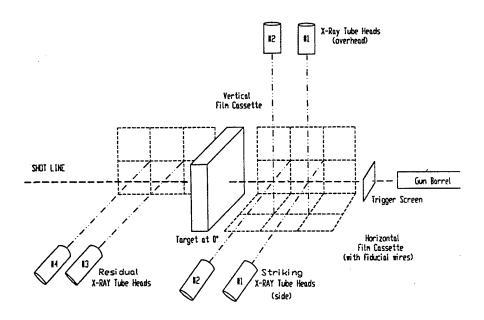


Figure 1. Experimental Setup.

The rod projectile used in the impact experiments had a mass of 65 g, a length (ℓ_0) of 78.2 mm, a diameter of 7.82 mm, and a hemispherical nose shape. The rods were composed of WA, to include tungsten (93%), nickel (4.91%), and iron (2.11%), and were manufactured by Teledyne Firth Sterling. The rods were fabricated using the nominal X-21 process with 25% swaging. Nominal material properties are as follows: density - 17.7 g/cm³, hardness - Rockwell C Scale 40.5–42.6, yield strength - 1.089–1.169 GPa, ultimate strength - 1.131–1.213 GPa, and elongation - 5.8–10.6 % [6, 7].

Targets included RHA (MIL-A-12560, Class 3 Steel) and titanium alloy, Ti-6/4 (6% aluminum and 4% vanadium). Two slightly different plate thicknesses were used for the RHA targets: 44.69 mm with an areal density of 350.8 kg/m² and a nominal 40 mm with a corresponding areal density of 307.3 kg/m². The Ti-6/4 target thickness was 70.21 mm, having an areal density of 312.4 kg/m².

Table 1 presents results of the impact experiments that were measured from the radiographs. It is to be noted that for both RHA and Ti-6/4 targets, higher rod-impact velocity resulted in greater rod residual length and velocity. The trend is not nearly so pronounced for residual rod length, but nonetheless, is consistent with expectations based on similar results by Zook and Frank [8] and Stilp [9]. The analysis of Grace [1] suggests that residual rod length is a slowly varying function of impact velocity, especially at relatively high velocities (over 1,500 m/s), when the rod overmatches the target substantially, as is the case here. On the other hand, residual rod velocity is expected to continue to increase at least linearly with impact velocity [1, 2, 3, 8], even at high velocity, and the present results show that trend.

Target channel profiles were measured from recovered targets after the shot and are given in Table 2. Figure 2 provides a schematic showing measurement locations relative to the target cross section. Two diameter measurements were recorded for each location to capture any possible noncircular shape of the channel. Mostly, there are only small differences in the gross channel features between the target materials beyond the presence of a spall ring at the Ti-6/4 target exit hole.

Only slight differences in entrance-hole diameters, resulting from cratering, and penetration-channel diameters were observed. Exit-hole diameters and spall-ring diameters appear to depend on

Table 1. Ballistic Results

Target	Pitch (°)	Yaw (°)	Striking Velocity (m/s)	Residual Rod Length (mm)	Residual Rod Velocity (m/s)
44.69-mm RHA 40.03-mm RHA	-0.25 -0.50 0.00	0.75 0.00 0.00	1,507 1,510 1,973	39.7 38.6 41.7	1,297 1,297 1,893
39.91-mm RHA 70.21-mm Ti-6/4	-0.75 -0.50	0.00	1,997 1,501	41.4 26.1	1,110
	-1.25 0.00 1.50	0.25 0.75 0.00	1,501 1,959 1,980	26.5 31.0 31.0	1,118 1,770 1,831

^a Lost second x-ray flash.

Table 2. Hole Measurements

	Entrance-Hole Diameter (mm)		Penetration Diamete		Exit-Hole Diameter (mm)	
Target	Min.	Max.	Min.	Max.	Min.	Max.
44.69-mm RHA 40.03-mm RHA 39.91-mm RHA	19 18 b b	20 20 b b	15 13 17 18	15 14 17 21	23 27 23/38 ^a 28/37 ^a	27 29 23/38 ^a 28/37 ^a
70.21-mm Ti-6/4	20 18 b b	27 22 b b	15 13 14 17	16 15 15 19	18/44 ^a 18/43 ^a 25/49 ^a 23/46 ^a	22/46 ^a 20/46 ^a 25/54 ^a 24/54 ^a

^a For measurements denoted xx/yy, yy is the spall ring diameter.

^b Obscured by Pusher Plate Impact.

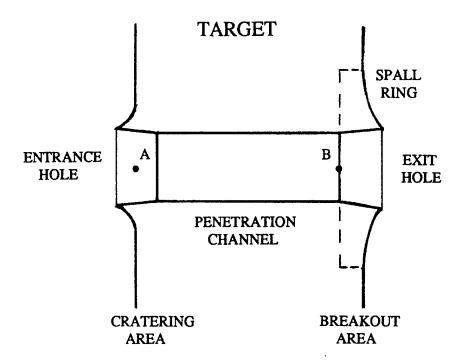


Figure 2. Target Cross Section Schematic, Where A Is the Original Surface, and B Is The Point Where Erosion Stops and Breakout Starts.

target material. Ti-6/4 exit-hole diameters averaged 20 mm and 24 mm in diameter for 1,500 m/s and 2,000 m/s, respectively, compared to the slightly larger 26.5-mm diameter averaged of the RHA steel plates. Penetration-channel diameter for the RHA demonstrated a slight increase with increased penetrator striking velocity. Ti-6/4's penetration channel may have increased with increased penetrator striking velocity, but the measurements are not conclusive. Plug thicknesses were obtained from flash x-ray coverage from behind the targets at 1,500 m/s. The average plug thickness was estimated to be 13 mm for the Ti-6/4 and 4 mm for the RHA. Plugs could not be separated from spall on the x-rays from the 2,000-m/s tests.

2.2 Microstructural and X-ray Diffraction Investigations. Both optical microscopy and transmission electron microscopy were used to study deformation of material near the channel wall. The midsection region of each target specimen was polished and etched for optical metallography.

For RHA, the etchant consisted of 2.5% nitol (2.5% nitric acid [HNO₃], 97.0% methanol [CH₃OH], and water [H₂O]) using etching times of roughly 35 s. The etchant for the Ti-6/4 target material consisted of 0.2-liter water, 5-ml hydrogen fluoride (HF), and 10-ml nitric acid, using etching times of approximately 17 s.

Specimens for transmission electron microscopy (TEM) were systematically extracted very close to the channel wall by using a stop-off lacquer on the channel surface and electropolishing from the rear. In addition, samples were sliced from representative sections at varying distances along the channel midsection, extending to distances of more than 20 mm from the channel surface. These slices were ground and polished to a thickness of about 100 μm, and 3-mm-diameter disks were punched from the slices. A Tenupol-3 electropolishing system was used to produce electron transparent thin sections utilizing the following etchants: for the RHA steel samples, 1.2 liters of methanol were mixed with 0.2 liters of perchloric acid (HClO₄), and electropolishing was performed at ~15° C; for Ti-6/4 alloy, 1.4 liters of methanol and 0.14 liters of hydrochloric acid (HCl) were employed at ~15° C. TEM analysis was done in a Hitachi H-8000 analytical TEM operated at 200-kV accelerating potential in the conventional TEM (CTEM) mode.

Figure 3(a) shows the longitudinal section of the RHA penetration channel. The corresponding optical (metallographic) views and TEM views of typical microstructures very near the channel surface are shown in Figure 3(b) and (d). Corresponding views at distances far removed from the channel are presented in Figure 3(c) and (e). There is extensive plastic deformation extending from the channel wall (shown in Figure 3[b]) that consists of gross grain elongation with very fine, elongated dislocation cell structures intermixed with carbides (shown in Figure 3[d]). Figure 3(e) illustrates the fact that the undisturbed RHA microstructure is a complex intermixing of ferrite grains, carbides, and a heavy dislocation structure. The average Vickers microhardness in that region is 0.32 GPa (using a 1-kgf load), while, in contrast, that taken at a distance of about 1 mm from the channel surface is about 0.4 GPa. Previously reported hardness profiles [10] taken outward from the channel surface indicate that the increased hardness and, therefore, the plastic zone due to the penetration process, extends radially to a distance of 70% of the channel radius or about 9.3 mm into the material beyond the channel surface.

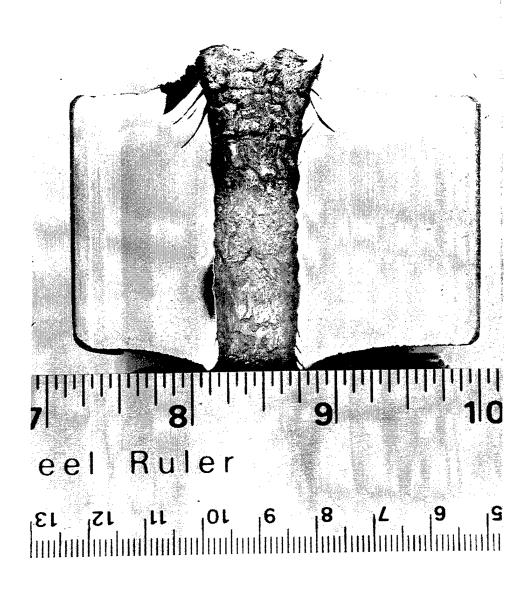
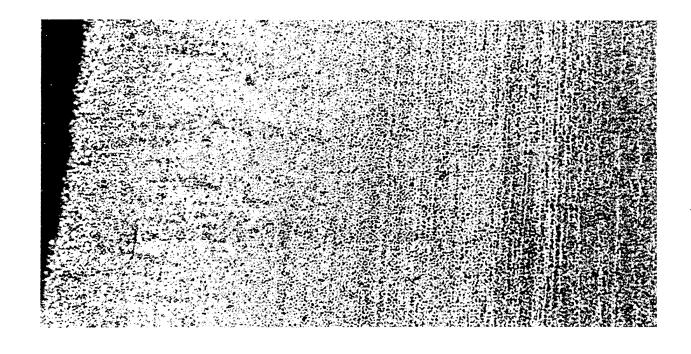


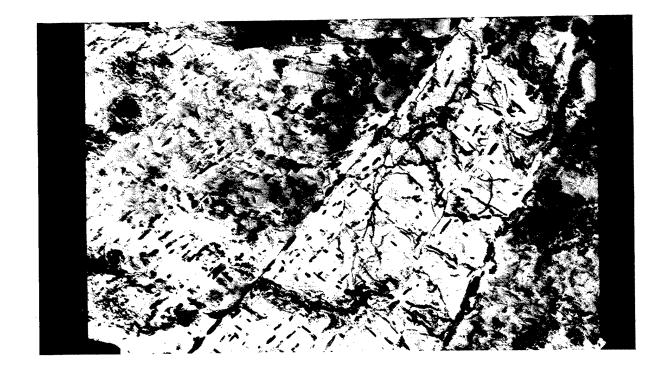
Figure 3. RHA Penetration-Channel-Related Microstructures. (a) RHA Penetration-Channel Half-Section Reference.



(b)

(c)

Figure 3. RHA Penetration-Channel-Related Microstructures (continued). (b) Light Microscope View of Microstructure Extending From the Penetration-Channel Surface at the Midsection. (c) Light Microscope View of the RHA Target Microstructure 20 mm From the Penetration-Channel Surface.



(d)

Figure 3. RHA Penetration-Channel-Related Microstructures (continued). (d) TEM Bright-Field Image of RHA Penetration-Channel-Related Microstructure Roughly 0.2 mm From the Penetration-Channel Surface. (e) TEM Bright-Field Image of the Initial RHA Target Microstructure Corresponding to (c).

(e)

Figure 4(a) shows the Ti-6/4 penetration-channel half section. TEM and optical microstructures close to the channel surface are shown in Figure 4(b) and (d), respectively. Their undisturbed counterparts at a distance far from the channel surface are shown in Figure 4(c) and (e). For Ti-6/4, no discernable differences in grain structure are seen in the optical microscopy at locations near to or far removed from the channel surface. This result is in stark contrast to that for RHA and suggests that no gross plastic deformation (grain elongation) exists in the Ti-6/4 material adjacent to the channel surface. However, TEM observations only 0.1 mm from the channel surface reveal (contrast Figure 4[c] and [e]) heavy dislocation structures within the 2-phase, α/β (hexagonal-close-packed [hcp]/body-centered-cubic [bcc]) regimes. It should be noted that grain elongation requires considerably more expended energy in deformation than does the generation of dislocations within the grains. The Ti-6/4 has a hardness profile [10] that shows an increase in hardness starting at the channel surface and continuing outward to about 40% of the channel radius or about 4 mm beyond the channel surface. There were no observations of shear bands associated with the channel midsection as shown for the Ti-6/4 target in Figure 4(b) and (d). Shear bands were also not prominent near the midsection of penetration channel in RHA as well.

Additional differences between the two target materials occur in the formation of the plugs and spall rings. The RHA exhibits signs of a higher degree of ductile failure in the form of material flow and tearing. In contrast, the Ti-6/4 target does not exhibit these gross features, but does resemble a more brittle failure. The conspicuous lack of plastic flow and tearing in Ti-6/4 was previously observed by Woodward, Baxter, and Scarlett [11] and Holt et al. [12] for plug formation at lower impact velocities. Further, the appearance of a spall ring in Ti-6/4 at 1,500-m/s impact velocity (and not in RHA) together with the previously mentioned observations suggest that Ti-6/4 is more susceptible to brittle failure under impact conditions than RHA.

As a check on any possible shifts in the α/β structure of the Ti-6/4 material during penetration, x-ray diffraction (Cu-K α) spectra were also compared between the initial material (far removed from the channel surface) and that within 1 mm of the channel surface. Results showed the β (bcc) phase to dominate, and there was no change in the material near the channel surface. All peak intensities

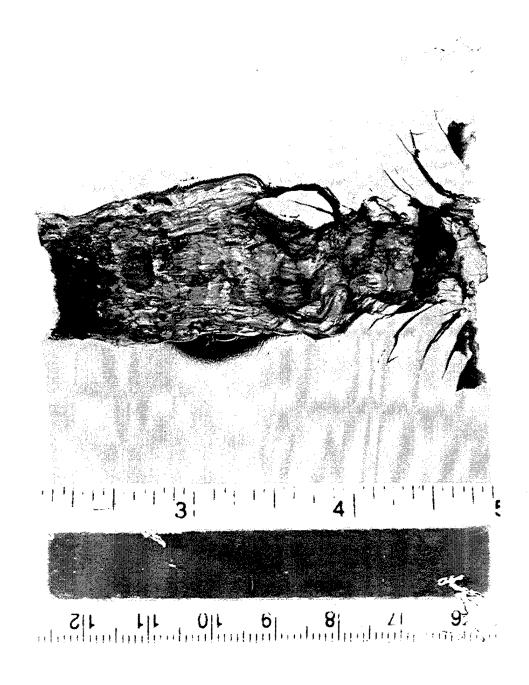
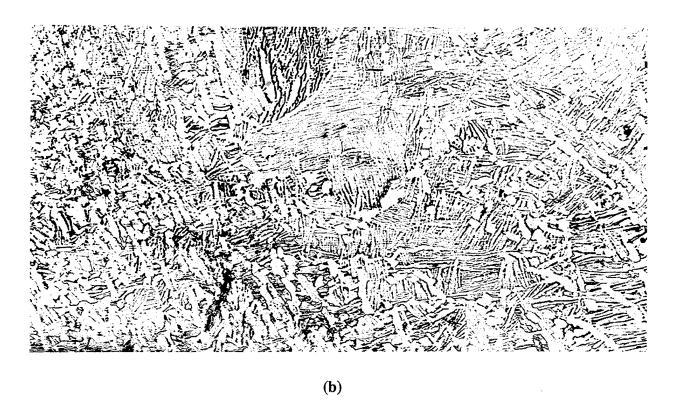


Figure 4. Ti-6/4 Penetration-Channel-Related Microstructures. (a) Ti-6/4 Penetration-Channel Half-Section Reference.



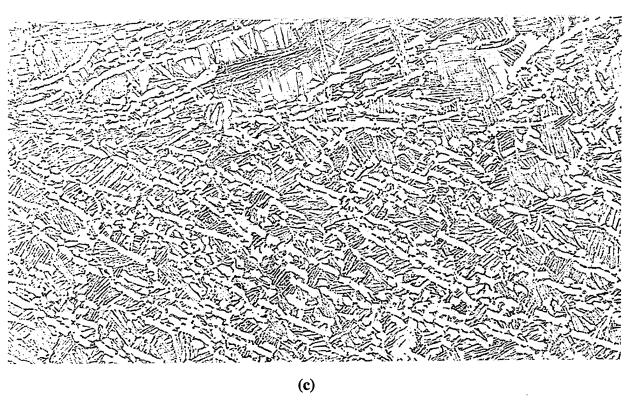
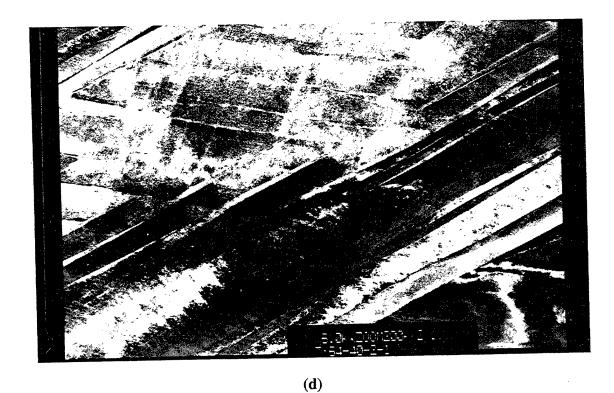


Figure 4. Ti-6/4 Penetration-Channel-Related Microstructures (continued). (b) Light Microscope View of Microstructure Extending From the Penetration-Channel Surface at the Midsection. (c) Light Microscope View of the Ti-6/4 Target Microstructure 20 mm From the Penetration-Channel Surface.



(e)

Figure 4. Ti-6/4 Penetration-Channel-Related Microstructures (continued). (d) TEM Bright-Field Image of Ti-6/4 Penetration-Channel-Related Microstructure Roughly 0.2 mm From the Penetration-Channel Surface. (e) TEM Bright-Field Image of the Initial Ti-6/4 Target Microstructure Corresponding to (c).

were enhanced uniformly for the material near the channel surface, which is consistent with the typical deformation-related dislocation density increases shown previously.

Since no gross plastic deformation was observed in Ti-6/4, nor an indication of a change in the distribution of phases within the Ti-6/4, it would appear that the penetration process must involve some other mechanism. Furthermore, since the Ti-6/4 has a tendency for brittle failure under impact as was discussed previously, it is plausible that the penetration involves extensive fracturing of the material. In Figure 3(a), a small chip can be seen clinging to the channel surface. Since the cavity has reasonable axial symmetry elsewhere along its length, it is of interest that no such material particle is present on the opposite side of the channel surface. The cracks about and within the chip suggest that the cavity may have been formed by excavation of fractured material.

3. Energy Analysis

Calculations of energy partitioning during sequential phases of penetration and perforation follow the analysis of Grace [1]. In this model, the penetration phase involves a simultaneous process of rod deceleration and erosion (reduction in its length) and target acceleration and erosion or plastic flow (reduction in its thickness). As such, the dynamics indicate that the rod length is not entirely consumed during penetration, nor is the target thickness. Thus, the analysis provides the unconsumed target thickness z_c , when target consumption (penetration) stops as

$$z_{c} = z_{0} \exp \left[-\frac{\rho_{t}}{2S_{t}} u_{0}^{2} \right], \qquad (1)$$

where z_0 is the target initial thickness, ρ_t is the target density, S_t is the target strength, and u_0 is the penetration rate. Grace [1] also provides a near-linear relationship between u_0 and striking velocity v_s . The consumed target thickness (penetration into the target) $P_c = z_0 - z_c$ corresponds to the distance between point A and point B in Figure 1. Also, the analysis provides residual rod length ℓ_r , based on its initial length ℓ_0 , and rod velocity v_e at point B and the acceleration of localized target material in front of the penetrator to a velocity u_0 as well. In general, the velocity of the remaining

rod at the end of the penetration phase will be greater than that of the target material (i.e. $v_e > u_0$). This condition promotes breakout. With masses and velocities defined by the momentum interactions contained within the theory, it is possible to calculate the associated KEs at the end of the penetration phase as inputs to the perforation phase, given an estimate for the cross-sectional area of the localized target material [2].

For the perforation phase, Grace [1] treats the subsequent rod-target interaction problem as inelastic so that after perforation, the residual rod and target plug velocities are equal. The energy balance under these circumstances is

$$(M_r + M_c)v_r^2 = M_r v_e^2 + M_c u_0^2 - 2E_f,$$
 (2)

where E_f is the energy required to fracture the remaining target material of thickness z_c , M_r is the residual rod mass, M_c is the residual target mass (assuming its area), and v_r is the residual rod velocity. Thus, KEs of the major masses are accounted for in the analysis. It is not necessary to account for the KE of the spall ring since the x-rays indicate that spall-ring velocity did not exceed one-tenth of the residual rod velocity even though its mass may be on the order of the plug mass.

The fracture model for the plug is essentially that of Woodward and Crouch [13] and Holt et al. [12], who determined the work required for plug separation (shear) from the surrounding target plate. In terms of Grace's notation, the fracture energy becomes

$$E_{f} = \frac{1}{2} \pi d_{s} S_{s} z_{c}^{2}, \qquad (3)$$

where S_s is the shear strength of the target material (taken here as $S_s = S_t/\sqrt{3}$, where S_t is a nominal strength for the target material), and d_s is the plug diameter. The energy required to fracture the spall ring is neglected for the cases of interest here. From the experiments it was noted that at lower velocities, the depth of the spall ring (as measured on the exit surface of the target) was about 5 mm for the Ti-6/4 while the plug thickness was 13 mm. Since fracture energy as given by equation (3) depends on the thickness squared, that for the spall ring is small (15%) relative to the fracture energy

of the plug. There was no observed spall ring on the RHA target impacted at 1,500 m/s. At high velocity, the fracture energy due to both plugging and spall-ring formation is relatively small compared to the high KEs of the rod and target masses. This notion follows from equations (1) and (3) since z_c gets exceedingly small as rod striking velocity and, therefore, penetration rate u_0 is increased.

Using the previous analysis, calculations were carried out for each of the rod-target combinations through a range of velocities that encompass the experimental conditions. The analysis used the average of all measured penetration-channel exit-hole diameters as the plug diameter. These diameters were 21 mm for RHA and 18 mm for Ti-6/4. Also, the following material properties were used: for RHA, $\rho_t = 7.85$ g/cm³, $S_t = 1.09$ GPa, $S_s = 0.629$ GPa, and $C_0 = 5,170$ m/s; for Ti-6/4, $\rho_t = 4.45$ g/cm³, $S_t = 1.14$ GPa, $S_s = 0.658$ GPa, and $C_0 = 6,070$ m/s. Properties for the tungsten alloy rods were taken to be $\rho_p = 17.3$ g/cm³ and $S_p = 1.51$ GPa, where S_p is the penetrator strength. Table 3 contains calculated results corresponding to the experimental impact conditions.

Table 3. Calculated Values

Target	Striking Velocity (m/s)	Residual Rod Length (mm)	Residual Rod Velocity (m/s)	Plug Mass (g)
44.69-mm RHA	1,507 1,510	39.91 38.96	1,271 1,275	12.508 12.360
40.03-mm RHA	1,973	48.64	1,904	1.252
39.91-mm RHA	1,997	48.92	1,931	1.095
70.21-mm Ti-6/4	1,501 1,959 1,980	27.93 34.23 34.52	1,096 1,769 1,797	20.758 5.243 4.871

Figure 5 presents a comparison of the calculations and experimental data for the residual rod length and velocity and plug thickness. Since the analytic results agreed well with the experimental results, they provided a valid basis for the energy analysis. The results of the energy calculations to include the KEs are presented in Table 4. The quantities of interest include initial rod KE before

impact, its KE after erosion and deceleration at the end of penetration into the target (first phase), the KE of accelerated target mass at the end of the first phase, the KE of the residual rod and target mass (plug) after perforation (second phase), and the fracture energy required to separate the plug from the target surrounds (fracture energy of breakout). The amount of KE dissipated during rod erosion is that lost by the rod during penetration (first phase) and is based on the difference between its initial KE before impact and that at the end of the penetration phase.

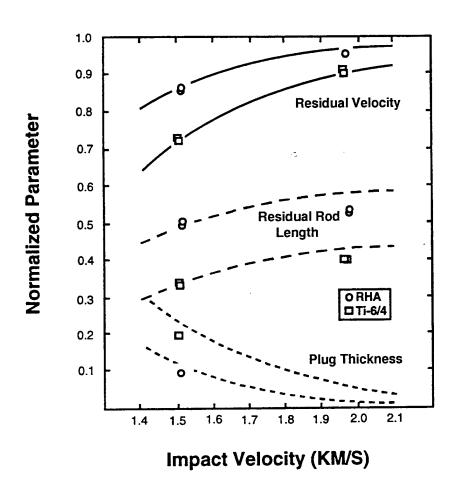


Figure 5. Comparison of Theoretical Calculations (Curves) and Experimental Data for RHA and Ti-6/4 to Include Residual Rod Velocity (v_r/v_s) , Residual Rod Length (ℓ_r/ℓ_0) , and Plug Thickness (z_c/z_0) .

The amount of KE dissipated from the rod by erosion and deceleration during penetration can appear in several forms to include KE of erosion products, work required to create the channel

Table 4. Energy Partition

Target	Striking Velocity (m/s)	Initial Rod KE (J)	KE Dissipated During Rod Erosion (J)	Plug KE Before Breakout (J)	Fracture Energy at Breakout (J)	Residual Rod KE (J)	Plug KE After Breakout (J)
44.69-mm RHA	1,507 1,510	75,206 75,506	41,629 42,543	6,271 6,141	177 172	27,129 26,650	10,102 10,046
40.03-mm RHA	1,973	128,909	53,224	1,487	2	74,196	2,269
39.91-mm RHA	1,997	132,064	53,951	1,357	1	76,755	2,041
70.21-mm Ti-6/4	1,501 1,951 1,980	74,609 126,050 129,826	53,475 76,434 78,542	4,863 4,406 4,206	2,154 137 119	14,117 45,073 46,905	12,467 8,204 7,865

cavity, acceleration of the target mass (plug), and heat. Table 4 indicates that the rod loses more KE to these factors when penetrating Ti-6/4 than it does in RHA plates. This result stems from the penetration path being longer in the Ti-6/4 while the target strengths (resistance to penetration) are similar. Also, since the plug thickness according to equation (1) depends on initial plate thickness, and the energy for plug fracture is greater by equation (3), considerably more plug fracture energy is expended in the Ti-6/4 target. Consequently, these two considerations provide a rationale for the greater efficiency of Ti-6/4 against these rods.

Table 5 provides the results of energy rate calculations and displaced target mass estimates. Time rates are based on the calculated time [1] between the initial contact of the penetrator at point A (Figure 2) and the time the penetrator stops eroding at point B (Figure 2). Length rates are based on length of rod eroded for the penetrator, the distance between the point of impact (point A, Figure 2), and the onset of breakout (point B, Figure 2) for the target. Mass rates are based on the mass eroded from the penetrator and the displaced target mass in the penetration channel (target).

When the amount of KE lost by the rod per unit path length in the target is considered, as shown in Table 5, the two target materials appear to be equivalent. Thus, the total energy expended is consistent with the work done on the target, which is proportional to the integrated effects of target

Table 5. Energy Rates and Displaced Mass

Target	KE Striking Dissipated		KE Dissipated Length (J/mm)		KE Dissipated Mass (J/g)		Displaced Target Mass in Penetration
Tangot	Velocity (m/s)	Time (J/µs)	Penetrator	Target	Penetrator	Target	Channel (g)
44.69-mm RHA	1,507 1,510	671 689	1,073 1,071	1,037 1,058	1,275 1,272	1,182 1,205	35.2 35.3
40.03-mm RHA	1,973	1,456	1,771	1,345	2,104	1,534	34.7
39.91-mm RHA	1,997	1,506	1,812	1,365	2,153	1,559	34.6
70.21-mm Ti-6/4	1,501 1,959 1,980	586 1,194 1,247	1,053 1,718 1,778	1,001 1,159 1,185	1,252 2,042 2,112	810 938 959	66.0 81.5 81.9

resistance times displacement. The displacement (path length) is greater for Ti-6/4. On the other hand, since it takes more time to penetrate the Ti-6/4 target, KE losses during penetration per unit time are similar for the two target materials.

Table 5 also shows that the penetrator mass and length based KE rates for Ti-6/4 are similar to that for RHA. Thus, the mechanism for creating cavities in Ti-6/4 must involve a process that is similar in energy requirements as plastic deformation in RHA. However, the microstructural observations validate that the Ti-6/4 did not exhibit extensive plastic deformation. Further, the present efforts have also shown no significant change in the α/β phase distribution; thus, this factor cannot contribute to energy absorption. Since Ti-6/4 is more susceptible to brittle fracture, it is possible that fracturing of the material initially located within the channel may occur during penetration. If this material could separate itself from the channel surface, possibly by shear, then this hypothesis is consistent with all of the observations and analyses conducted in this study. Thus, there is reason to pursue an energy-consuming mechanism during the creation of the channel in Ti-6/4 that involves extensive fracturing of Ti-6/4 material and its excavation during the penetration process.

4. Summary of Results

The problem of target penetration and perforation has been addressed using multiple approaches to include (1) an energy analysis of the penetration and perforation processes, (2) specifically designed long rod-impact tests using tungsten alloy rods at 1,500 and 2,000 m/s against RHA and Ti-6/4 plates of equal areal density, and (3) metallurgical investigations of material adjacent to penetration channels created by impact. The experimental work and the analysis showed that Ti-6/4 is a substantially more mass-efficient target material than RHA. The dynamics contained within the analysis showed that the increased efficiency of the Ti-6/4 was due to higher energy dissipation by the rod in that target material. While the Ti-6/4 target strength was similar to RHA, penetration paths (amount of target thickness) and, therefore, the forces times displacements were considerably larger for Ti-6/4 plates of equal areal density. Plug thicknesses were larger for the Ti-6/4, so the additional energy required for plug separation also contributed.

The experiments and associated calculations indicated that length-based KE dissipation rates were similar in the formation of the penetration channels in both Ti-6/4 and RHA. However, the metallurgical observations of the channel walls showed a decisive difference in the process by which penetration takes place. The grain elongation and deformation structures observed show that RHA follows the classical cavity expansion related to plastic flow about the rod during penetration. On the other hand, the Ti-6/4 did not exhibit these gross plastic deformation features. Thus, since the Ti-6/4 has known tendencies for brittle failure, and since some evidence of brittle failure was observed within the penetration channel, it is plausible that Ti-6/4 undergoes extensive fracturing with possible excavation of the fractured material during penetration. Such a process would not require large amounts of plastic flow, but rather much more energy expended in fracture, and this view is consistent with the collective experimental observations and analyses conducted in this investigation.

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